

Fluid mechanics and granular matter

At ENSIMAG, bat. H, room H102

Monday 14h-17h

Lectures in two parts:

A. Fluid mechanics

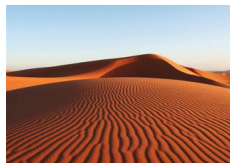
Pierre Saramito, $6 \times 3\text{h}$

- ▶ 23, 30 September 2024
- ▶ 7, 14, 21 October 2024
- ▶ 4 November 2024

B. Granular matter

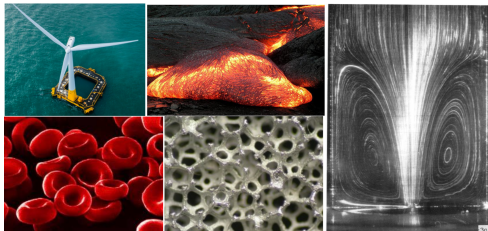
Didier Bresch, $6 \times 3\text{h}$

- ▶ 18, 25 November 2024
- ▶ 2, 9, 16 December 2024
- ▶ 6 January 2025



Aim: computer science \implies natural hazards, health, industry

A. Fluid mechanics

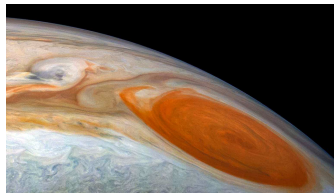
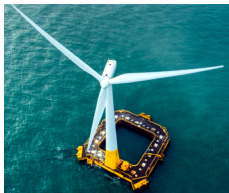


Fives chapters:

1. Navier-Stokes equations : 3h
2. Quasi-Newtonian fluids : 3h
3. Visco-plasticity : 4h30
4. Visco-elasticity : 6h
5. Elasto-visco-plasticity : 1h30

1. Navier-Stokes equations

for simple fluids only : air & water



(P): find \mathbf{v} and p such that

$$\begin{cases} \rho(\partial_t \mathbf{v} + (\mathbf{v} \cdot \nabla) \mathbf{v}) - \eta \Delta \mathbf{v} + \nabla p = \rho \mathbf{g} \\ \operatorname{div} \mathbf{v} = 0 \end{cases}$$

2. Quasi-Newtonian fluids: η non-constant

3. Visco-plasticity

= for mushy fluids : paste, mud

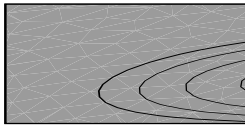
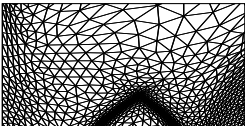
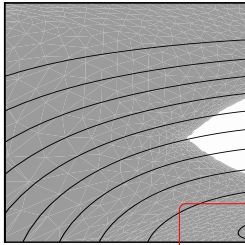
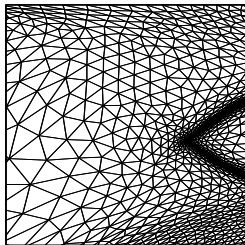
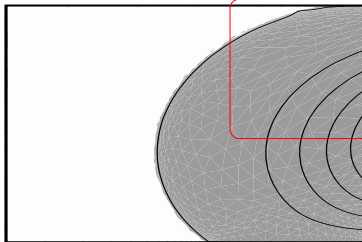
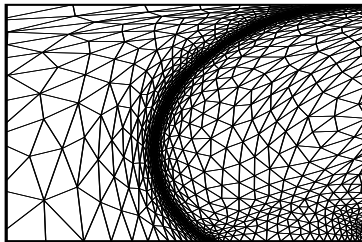


$$(P) : \min_{v \in H_0^1(\Omega)} J(v)$$

$$J(v) = \int_{\Omega} |\nabla v|^2 dx + \sigma_0 \int_{\Omega} |\nabla v| dx - \int_{\Omega} f v dx$$

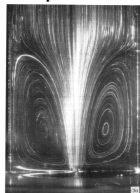
min : J is non-differentiable : $j(x) = x^2 + \sigma_0|x| - fx$

- optimization & convex analysis
- automatic adaptive mesh



4. Visco-elasticity

= for suspensions of long elastic molecules : polymers, biology



(P): find $(\boldsymbol{\tau}, \mathbf{v}, p)$ such that :

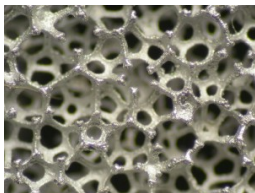
$$\begin{cases} \lambda \dot{\boldsymbol{\tau}} + \boldsymbol{\tau} - \nabla \mathbf{v} - (\nabla \mathbf{v})^T & = 0 \\ \operatorname{div} \boldsymbol{\tau} + \varepsilon \Delta \mathbf{v} - \nabla p & = \mathbf{f} \\ \operatorname{div} \mathbf{v} & = 0 \end{cases}$$

$$\dot{\boldsymbol{\tau}} = \partial_t \boldsymbol{\tau} + (\mathbf{v} \cdot \nabla) \boldsymbol{\tau} + \nabla \mathbf{v} \boldsymbol{\tau} - \boldsymbol{\tau} \nabla \mathbf{v}$$

- upwind schemes : discontinuous Galerkin methods
- mixed finite elements : inf-sup condition

5. Elasto-visco-plasticity

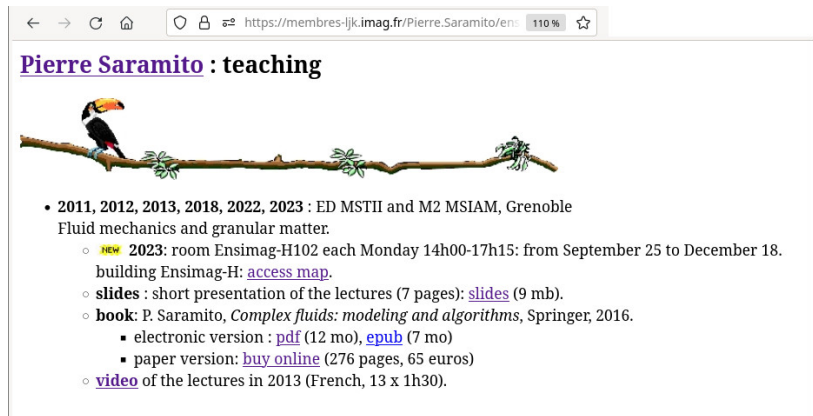
= for soft grains : biology, liquid foam



$$\lambda \dot{\boldsymbol{\tau}} + \max \left(0, 1 - \frac{\sigma_0}{|\boldsymbol{\tau}|} \right) \boldsymbol{\tau} = \boldsymbol{\nabla} \mathbf{v} + (\boldsymbol{\nabla} \mathbf{v})^T$$

⇒ combination of previous problems with (λ, σ_0)

Teaching material: on my web page

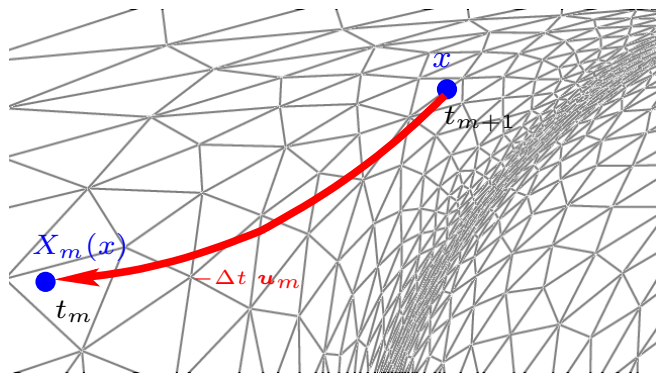


The screenshot shows a web browser window with the address bar containing the URL <https://membres-ljk.imag.fr/Pierre.Saramito/ens>. The page title is "Pierre Saramito : teaching". Below the title is an illustration of a toucan perched on a branch. The main content is a list of teaching activities:

- **2011, 2012, 2013, 2018, 2022, 2023** : ED MSTII and M2 MSIAM, Grenoble Fluid mechanics and granular matter.
 - **NEW 2023**: room Ensimag-H102 each Monday 14h00-17h15: from September 25 to December 18. building Ensimag-H: [access map](#).
 - **slides** : short presentation of the lectures (7 pages): [slides](#) (9 mb).
 - **book**: P. Saramito, *Complex fluids: modeling and algorithms*, Springer, 2016.
 - electronic version : [pdf](#) (12 mo), [epub](#) (7 mo)
 - paper version: [buy online](#) (276 pages, 65 euros)
 - **video** of the lectures in 2013 (French, 13 x 1h30).

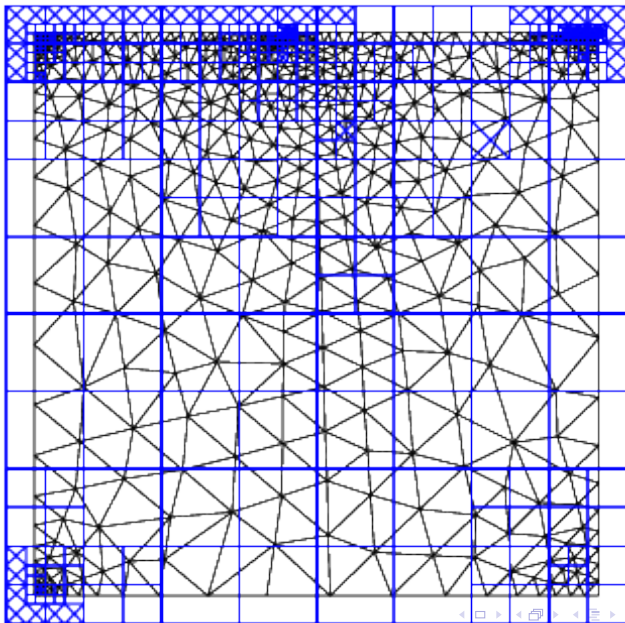
Evaluation of characteristics

$$X_m(x) \approx x - \Delta t \mathbf{u}_m(x)$$

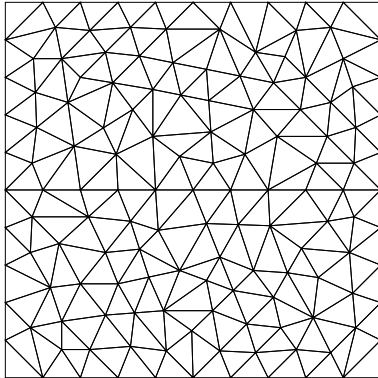


\implies interpolation : $\mathbf{u}_m(X_m(x))$

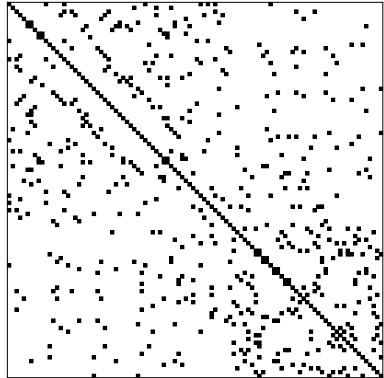
Searching: quadtree



Mesh & sparse matrix



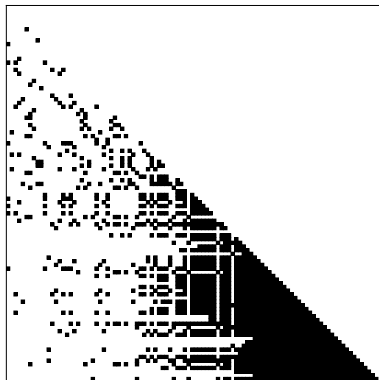
\mathcal{T}_h



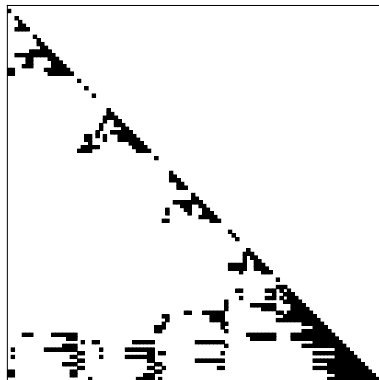
A

Sparse factorization $A = LDL^T$: renumbering

natural



optimized



Implementation with Rheolef (C++ library)

```
int main (int argc, char** argv) {  
  geo omega (argv[1]);  
  space Qh (omega, "P1");  
  space Xh (omega, "P2", "vector");  
  Xh.block ("boundary");  
  trial u(Xh), p(Qh); test v(Xh), q(Qh);  
  form a = integrate (2*ddot(D(u),D(v)));  
  form b = integrate (-q*div(u));  
  form m = integrate (p*q);  
  field uh (Xh, 0), ph (Qh, 0);  
  uh[1]["top"] = 0;  
  solver_abtb stokes (a.uu(), b.uu(), m.uu());  
  stokes.solve(-a.ub()*uh.b(), -b.ub()*uh.b(),  
    uh.set_u(), ph.set_u());  
  cout << catchmark("u") << uh  
    << catchmark("p") << ph;  
}
```

⇒ 15 lines of code

soit $\Omega \subset \mathbb{R}^d$, $d = 1, 2, 3$

$$Q_h = \{q \in L^2(\Omega); q|_K \in P_1, \forall K \in \mathcal{T}_h\}$$

$$X_h = \{\mathbf{v} \in H^1(\Omega)^d; \mathbf{v}|_K \in (P_2)^d, \forall K \in \mathcal{T}_h\}$$

$\forall \mathbf{u}, \mathbf{v}$, et p, q , définissons :

$$a(\mathbf{u}, \mathbf{v}) = \int_{\Omega} 2D(\mathbf{u}) : D(\mathbf{v})$$

$$b(\mathbf{u}, q) = - \int_{\Omega} q \operatorname{div} \mathbf{u}$$

$$m(p, q) = \int_{\Omega} p q$$

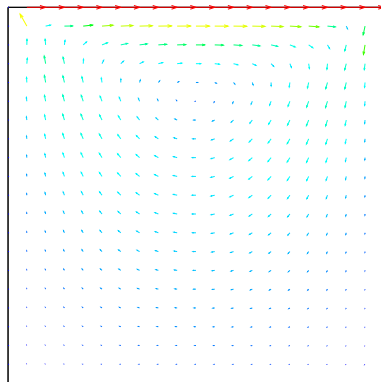
trouver $\mathbf{u} \in X_h$, $\mathbf{u}_h = 0$ sur $\partial\Omega$, et $p_h \in Q_h$ tels que

$$a(\mathbf{u}_h, \mathbf{v}_h) + b(\mathbf{v}_h, p) = 0, \quad \forall \mathbf{v}_h \in X_h, \mathbf{v}_h = 0 \text{ sur } \partial\Omega$$

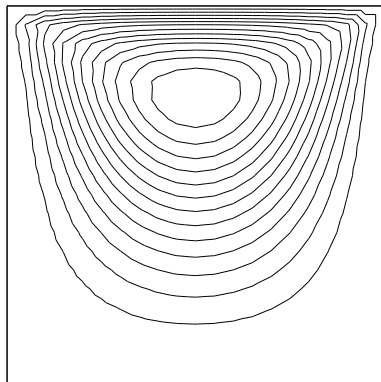
$$b(\mathbf{u}_h, q) = 0, \quad \forall q_h \in Q_h$$

Stokes in the driven cavity

velocity

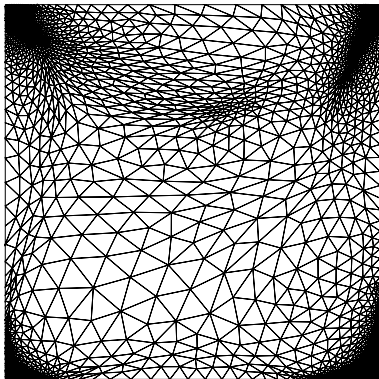


stream function

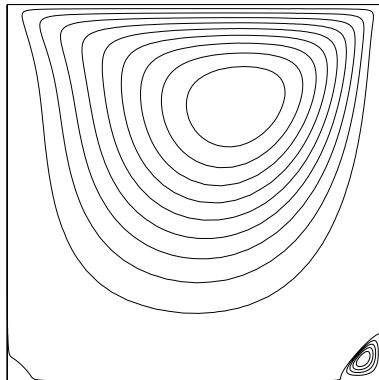


Navier-Stokes $Re = 100$: 4804 elements

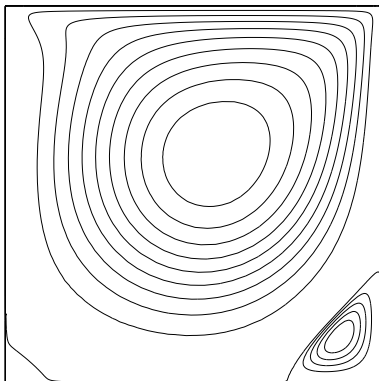
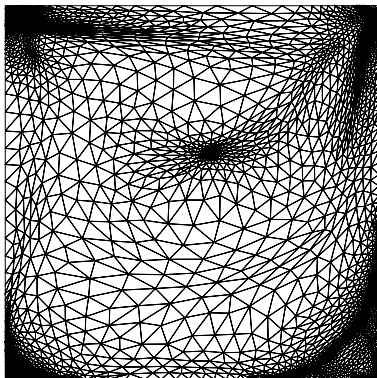
mesh



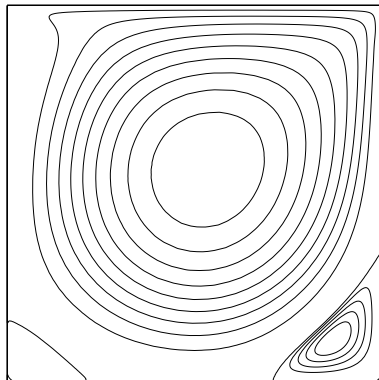
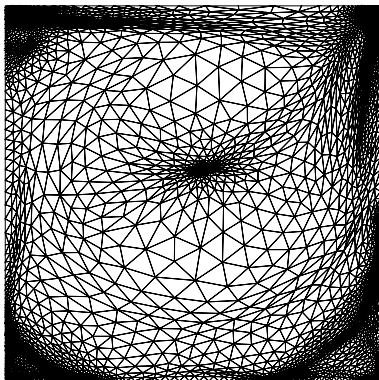
stream function



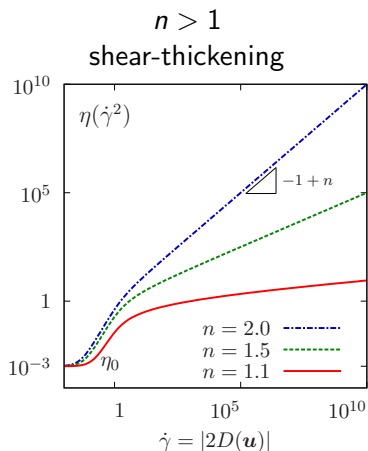
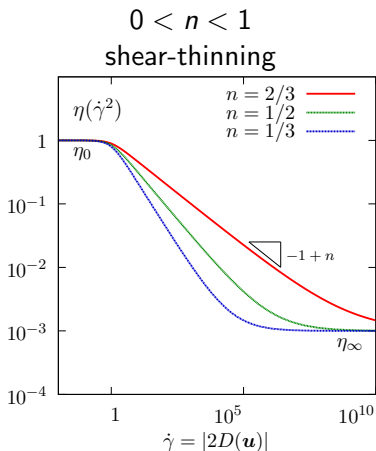
Navier-Stokes $Re = 400$: 5233 elements



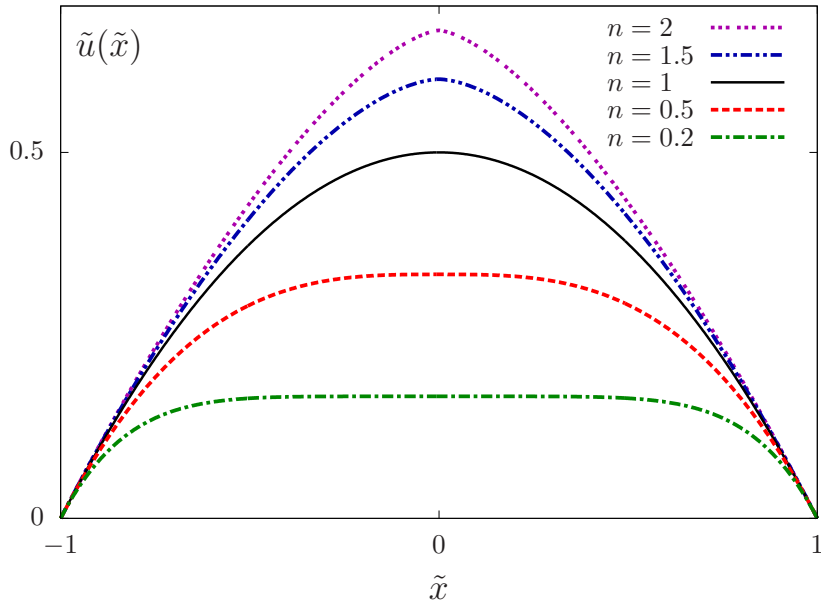
Navier-Stokes $Re = 1000$: 5873 elements



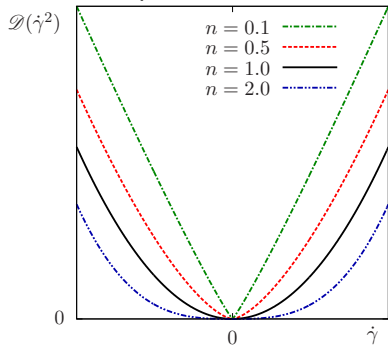
Viscosity : Carreau's law



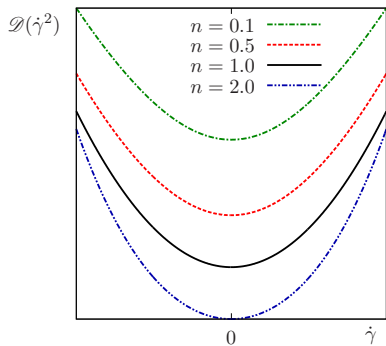
Poiseuille flow with power law



power law

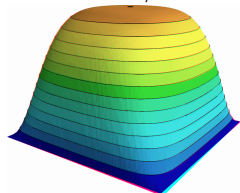


Carreau's law

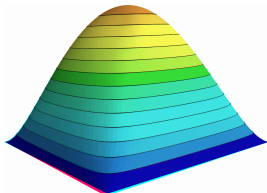


Poiseuille flow in a square pipe section

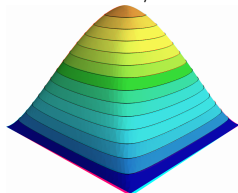
$$n = 1/2$$



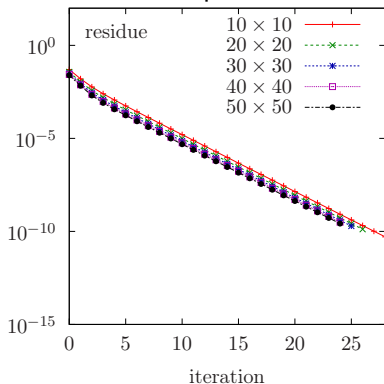
$$n = 1$$



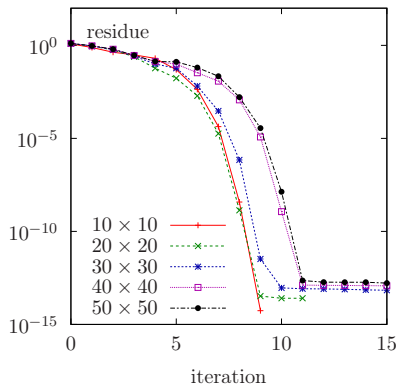
$$n = 3/2$$



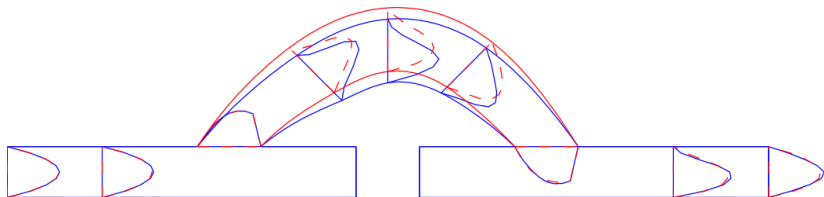
Fixed point



Newton

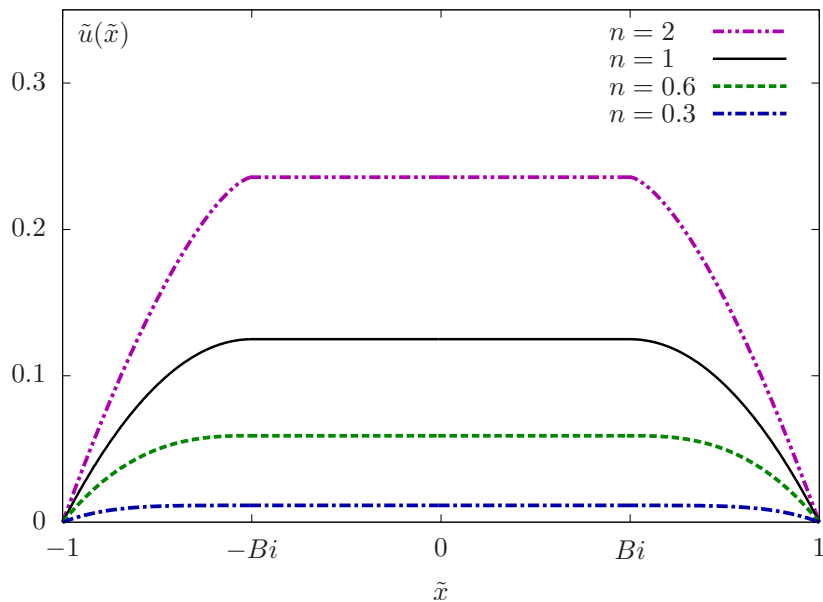


Example 3: arterial graft shape optimization

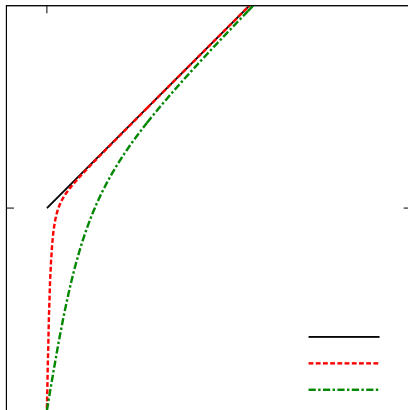


comparison: **Newtonian** and **Carreau's law** (Abraham *et al.* 2005)

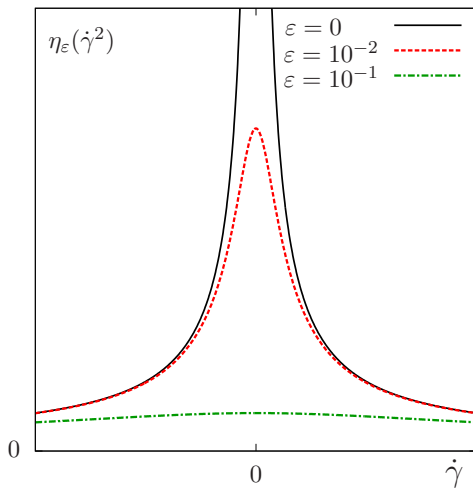
Poiseuille flow with viscoplastic fluid



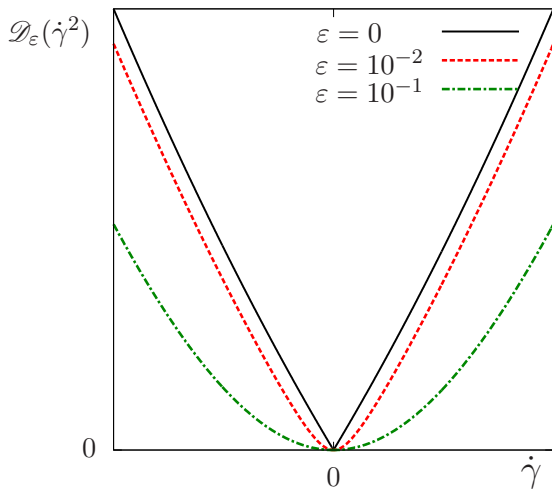
Regularization: stress vs shear



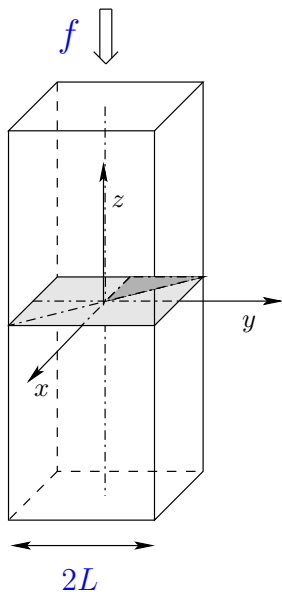
Regularization: viscosity



Dissipation: convexity



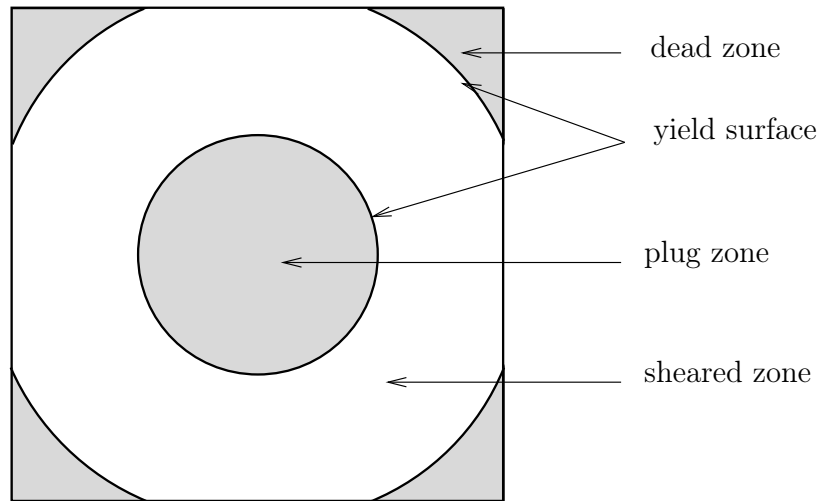
Poiseuille in a square section



dimensionless number

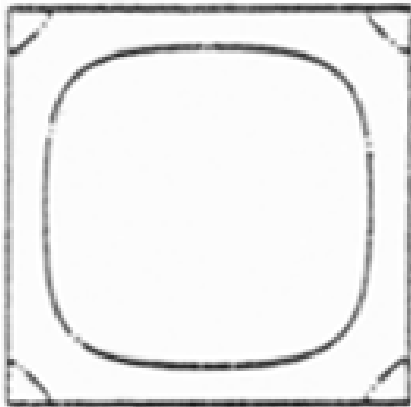
$$Bi = \frac{\sigma_0}{L f}$$

Terminology

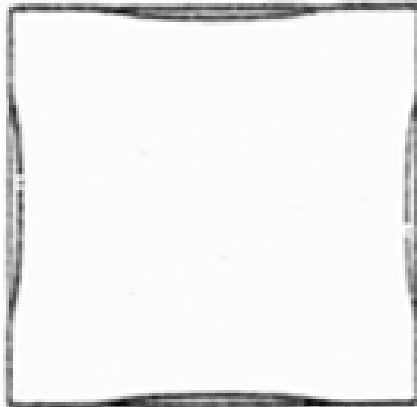


Regularization: Taylor & Wilson (1997)

$Bi = 0.8$



$Bi = 1.0$



⇒ badly shaped rigid zones...

Non-convergence $\varepsilon \rightarrow 0$? Wang (1999)

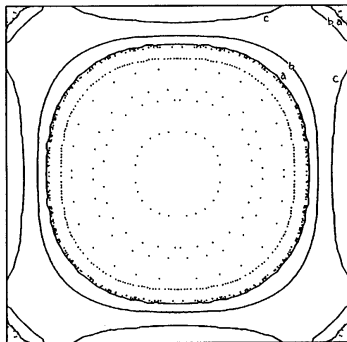
Taylor & Wilson

$$\varepsilon = 10^{-3}$$



Wang

$$\varepsilon = 10^{-13}$$



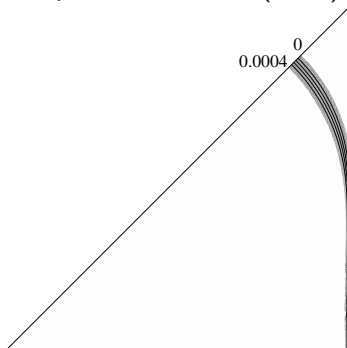
	$ 2D(u_\varepsilon) $
(c)	10^{-2}
(b)	10^{-3}
(a)	8×10^{-7}

Augmented Lagrangian algorithm

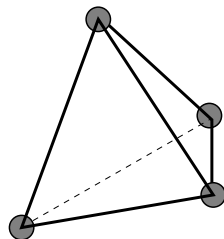
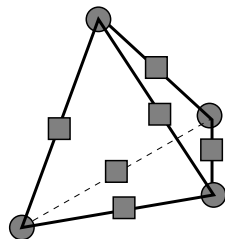
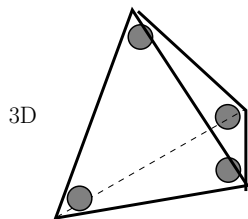
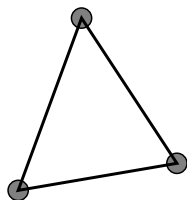
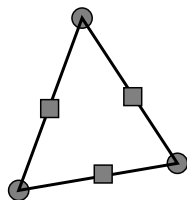
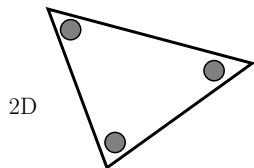
Taylor & Wilson (1997)



Roquet & Saramito (2001)



Viscoplastic fluids: space approximation

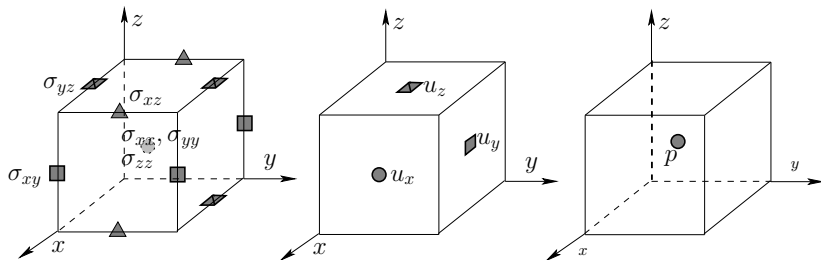
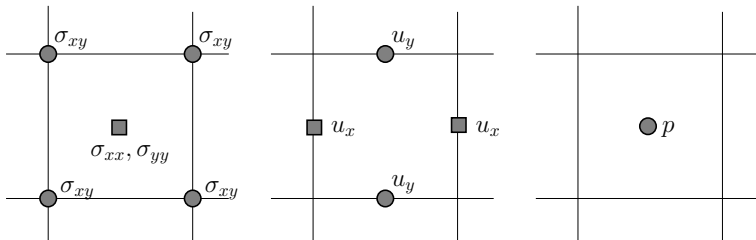


σ and γ : $P_1 - C^{-1}$

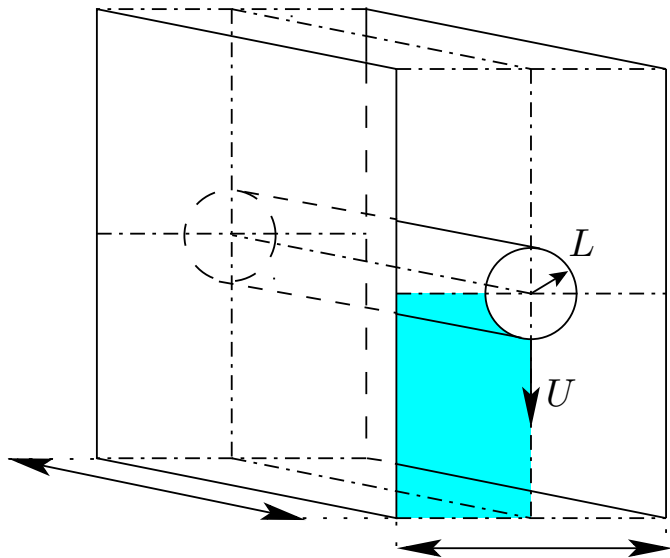
\mathbf{u} : $P_2 - C^0$

p : $P_1 - C^0$

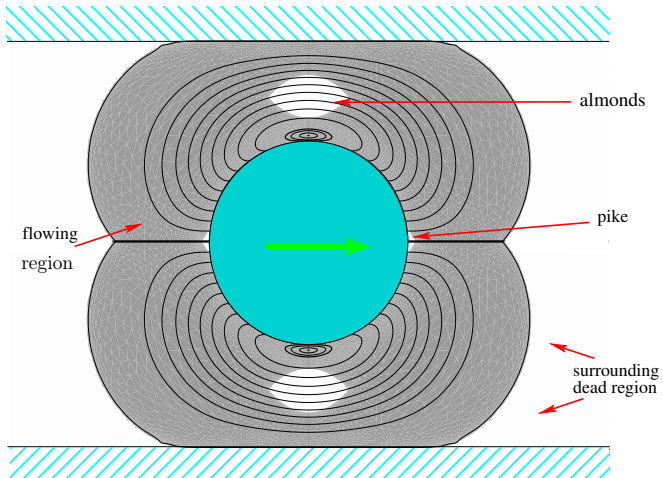
Finite differences



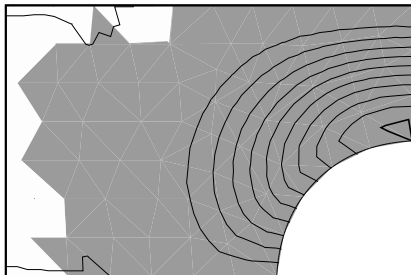
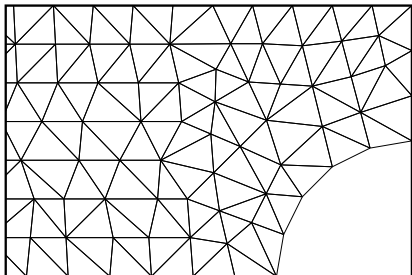
3.9 Example: flow around an obstacle



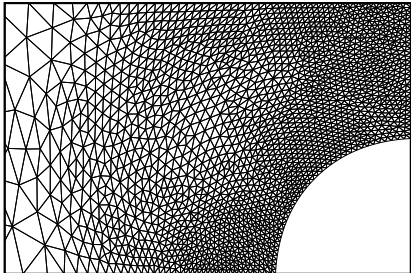
Terminology



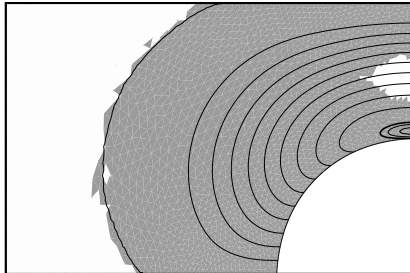
$$\text{card}(\mathcal{T}_h) = 539$$



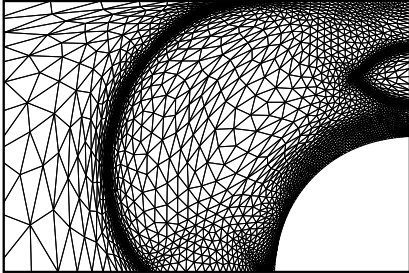
$$\text{card}(\mathcal{T}_h) = 15\,466$$



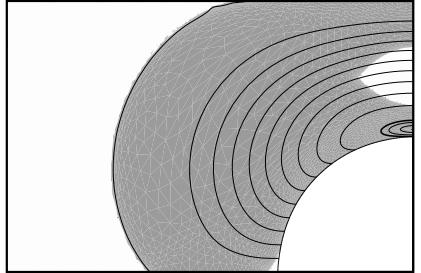
After mesh adaptation

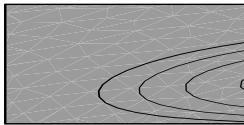
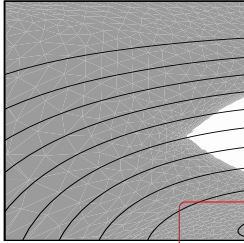
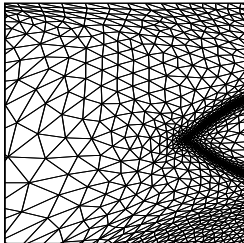
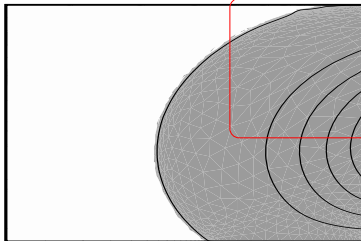
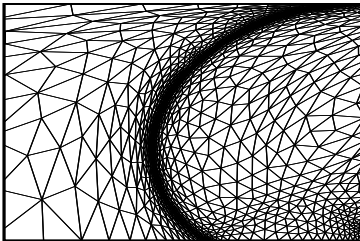


$$\text{card}(\mathcal{T}_h) = 41\,955$$

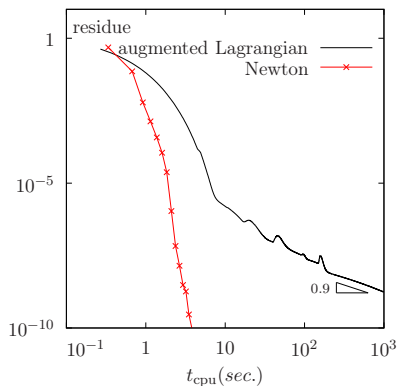
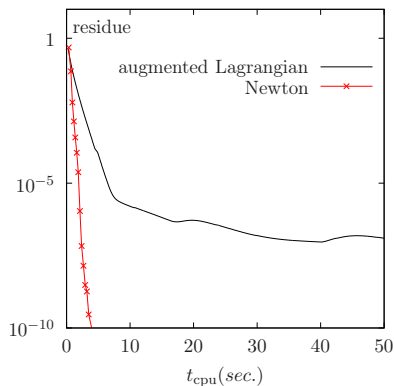


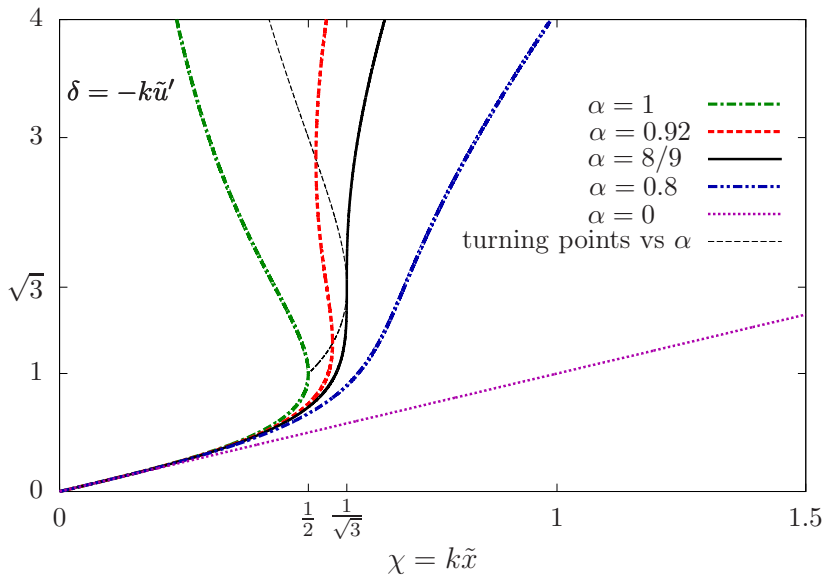
After 9 mesh adaptations



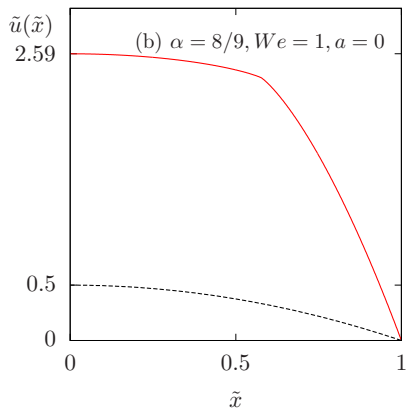
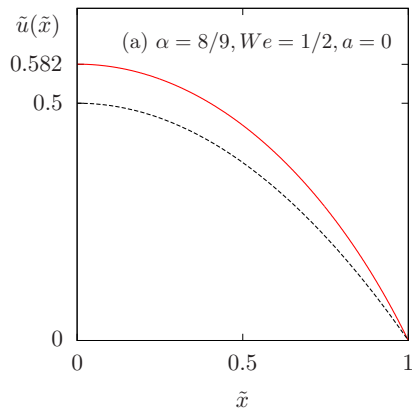


Viscoplastic fluid: augmented Lagrangian vs Newton

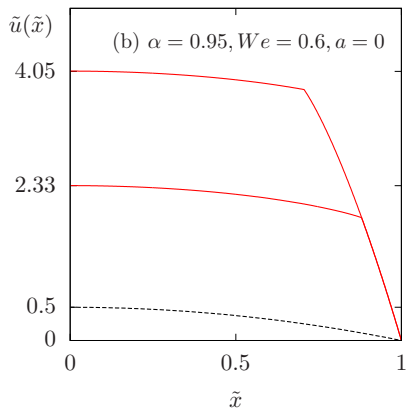
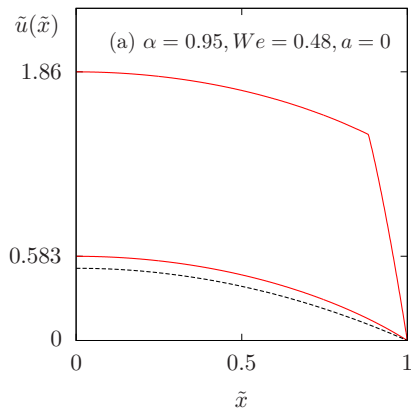




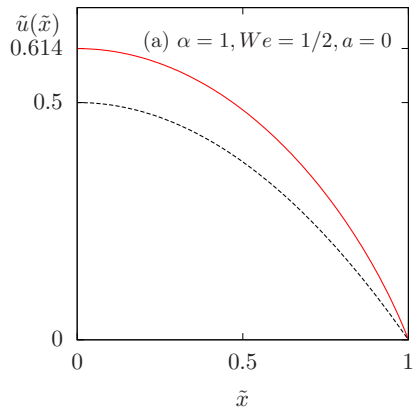
Poiseuille: Oldroyd $\alpha = 8/9$



Poiseuille: Oldroyd $\alpha = 0.95$



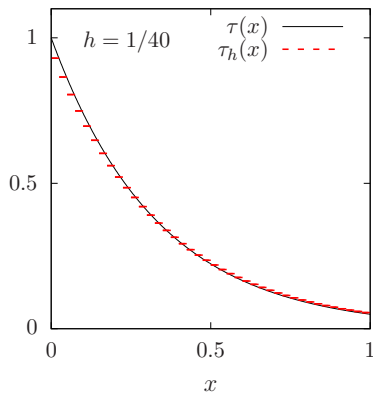
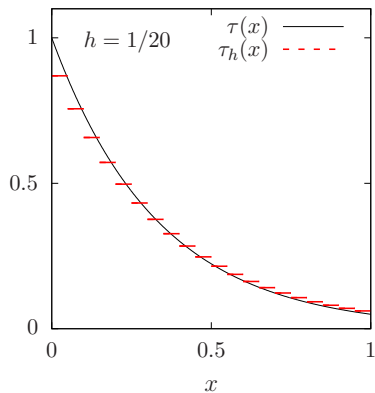
Poiseuille: Maxwell $\alpha = 1$



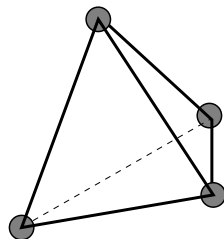
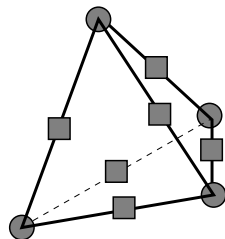
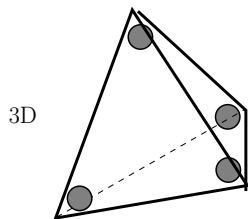
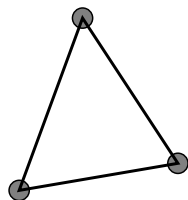
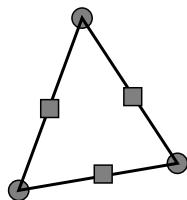
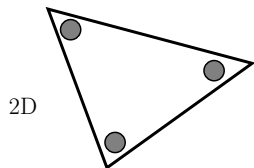
C++ code with Rheolef: scalar transport

```
1 #include "rheolef.h"
2 using namespace rheolef;
3 using namespace std;
4 int main(int argc, char**argv) {
5     environment rheolef (argc, argv);
6     geo omega (argv[1]);
7     space Xh (omega, argv[2]);
8     Float alpha = (argc > 3) ? atof(argv[3]) : 1;
9     Float sigma = (argc > 4) ? atof(argv[4]) : 3;
10    point u (1,0,0);
11    trial phi (Xh); test psi (Xh);
12    form ah = integrate (dot(u,grad_h(phi))*psi + sigma*phi*psi)
13              + integrate ("boundary", max(0, -dot(u,normal()))*phi*psi)
14              + integrate ("internal_sides",
15                          - dot(u,normal())*jump(phi)*average(psi)
16                          + 0.5*alpha*abs(dot(u,normal()))*jump(phi)*jump(psi));
17    field lh = integrate ("boundary", max(0, -dot(u,normal()))*psi);
18    solver sah (ah.uu());
19    field phi_h(Xh);
20    phi_h.set_u() = sah.solve(lh.u());
21    dout << catchmark("sigma") << sigma << endl
22          << catchmark("phi") << phi_h;
23 }
```

Example: scalar transport



Viscoelastic fluids: space approximation

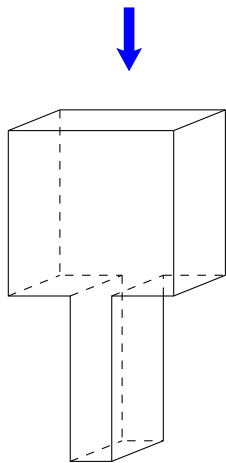


σ and γ : $P_1 - C^{-1}$

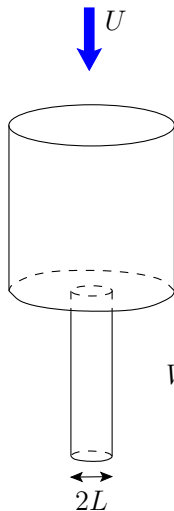
\mathbf{u} : $P_2 - C^0$

p : $P_1 - C^0$

Example: abrupt contraction



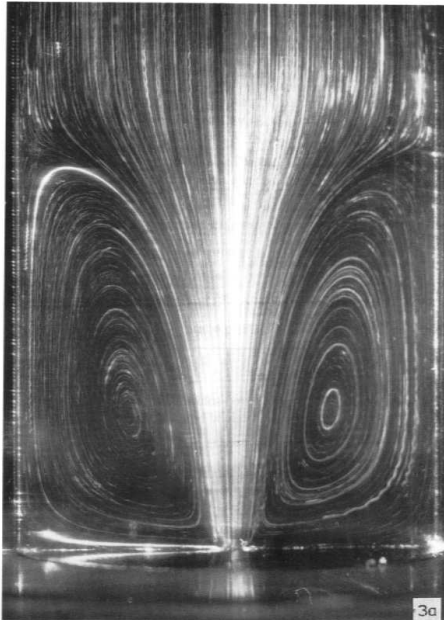
planar



axisymmetric

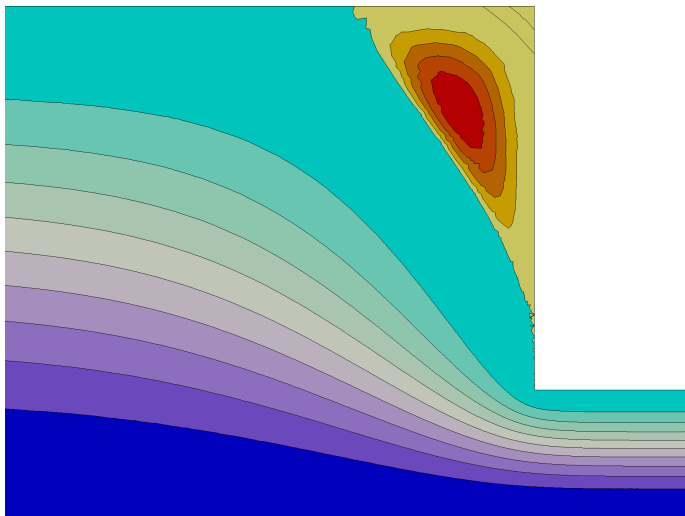
$$We = \frac{\lambda U}{L}$$

Polymer solution flow



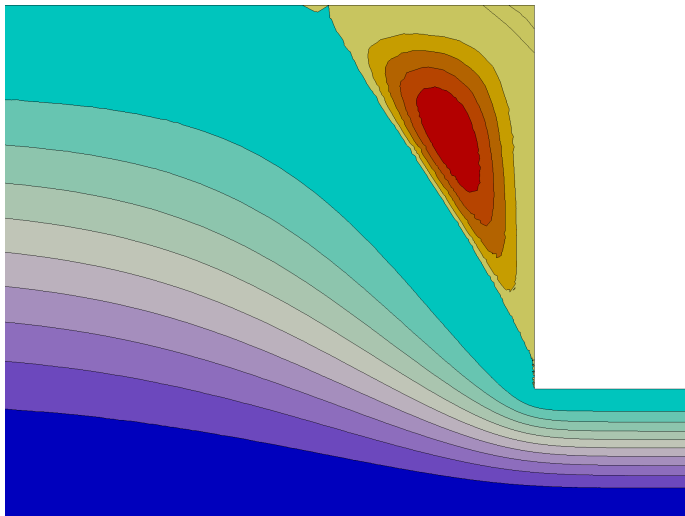
Oldroyd model: $We = \lambda U/L = 0.1$

Computations with Rheolef



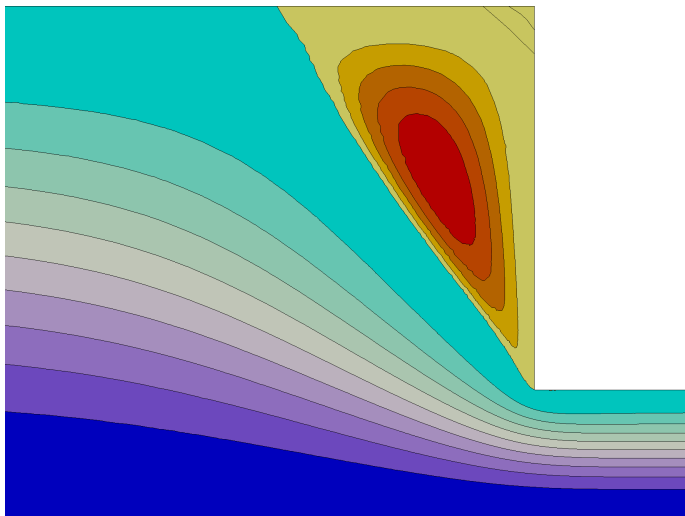
Oldroyd model: $We = \lambda U/L = 0.3$

Computations with Rheolef



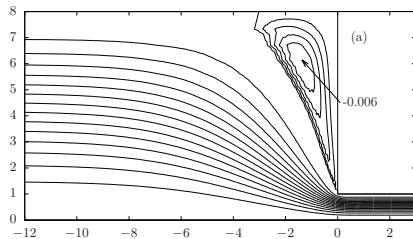
Oldroyd model: $We = \lambda U/L = 0.7$

Computations with Rheolef

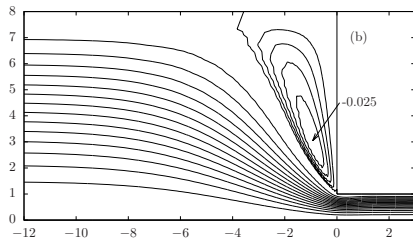


Phan-Thien and Tanner model

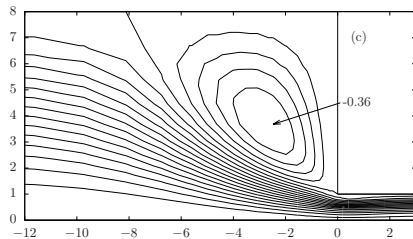
$We = 2$



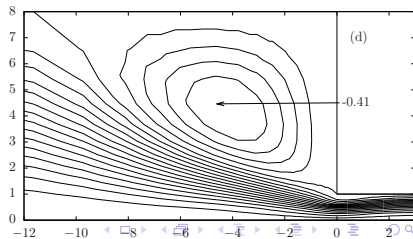
$We = 4.6$



$We = 14$



$We = 69$



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14h	Octave Crespel	Quelques résultats nouveaux sur les méthodes de projection
14h20	Basile Dubois Bonnaire	Shape optimization in steady blood flow
14h40	Jules Treton	Submerged jet shearing of viscoplastic sludge
15h	Hanne Christiaensen	Numerical modeling of shallow viscoplastic fluids
15h20	Sara Avesani	Viscoplastic flows with stick-slip at the wall
15h40	Yannick Godammer	Constitutive laws for the matrix-logarithm of the conformation tensor
16h	Guilia Quarta	Stationary Oldroyd model with diffusive stress
16h20	Yann Vincent	A new operator splitting algorithm for elastoviscoplastic flows
16h40	Antoine Sese	Particle migration in channel flow of an elastoviscoplastic fluid